

## PHYSICS 715

### Problem Set 5

Due Friday, March 3, 2006

Reading: Landau and Lifshitz, Secs. 32, 74–79

#### LD 14: The virial theorem and the equation of state of a neutral plasma.

A neutral plasma consists of  $N$  electrons with mass  $m_e$  and charge  $-e$ , labelled  $i = 1, \dots, N$ , and  $N$  ions with mass  $M$  and charge  $+e$ , labelled  $i = N + 1, \dots, 2N$ , all confined at temperature  $T$  in an insulating box of volume  $V$ . The internal Hamiltonian for the system is

$$H = \sum_{i=1}^N \frac{\mathbf{p}_i^2}{2m_e} + \sum_{i=N+1}^{2N} \frac{\mathbf{p}_i^2}{2M} + \sum_{\substack{i,j=1 \\ i < j}}^{2N} \frac{e_i e_j}{|\mathbf{x}_i - \mathbf{x}_j|}.$$

(Neutral species and more charged species could also be included.) Show that the equation of state for the plasma can be written as

$$P = nkT \left( 1 + \frac{E_{\text{Coulomb}}}{3kT} \right), \quad \text{where} \quad E_{\text{Coulomb}} = \frac{1}{2N} \left\langle \sum_{i < j} \frac{e_i e_j}{|\mathbf{x}_i - \mathbf{x}_j|} \right\rangle$$

is the average Coulomb interaction energy per particle in the exact distribution, and  $n = 2N/V$  is the total number density of the particles. [Hint: use the general form of the virial theorem.] Obtain an explicit expression for  $E_{\text{Coulomb}}$  in terms of electron-electron, ion-ion, and electron-ion interaction integrals, with all equivalent terms counted and combined as much as possible. Is the pressure increased or reduced by the interactions? Explain.

#### LD 15: The grand partition function $\mathcal{Y}(T, P, N)$ , density fluctuations, and a mean-field description of critical opalescence.

We can define a grand partition function  $\mathcal{Y}$  for systems in which  $T$  and  $N$  are fixed, but in which  $V$  can vary (*e.g.*, a small sample of a larger volume of gas) in a way similar to that in which we introduced the grand partition function  $\mathcal{Z}$ . We calculate the canonical partition function  $Z(T, V, N)$  for a given  $V$ , multiply by a factor  $y^V = e^{-\beta PV}$ , sum (integrate) over all possible volumes, and define

$$\mathcal{Y}(T, P, N) = \frac{1}{V_0} \int_0^\infty e^{-\beta PV} Z(T, V, N) dV,$$

where  $V_0$  is an arbitrary small volume included for dimensional reasons.

- (a) It may be shown that  $G(T, P, N) = -kT \ln \mathcal{Y}$  is the Gibbs free energy of the system,  $G = E - TS + PV$ . Show that this identification is correct for the ideal monotonic gas by calculating  $\mathcal{Y}$  and  $G$  explicitly, and showing that the appropriate derivatives of  $G$  give the correct values for the specific volume  $v = V/N$ , the entropy, and the chemical potential.

- (b) Density fluctuations lead to the scattering of light in fluids, *e.g.*, the blueness of the sky or critical opalescence near a liquid-gas phase transition, and a corresponding attenuation of the intensity of a beam of light with an attenuation coefficient

$$\alpha = \frac{8\pi^3}{3} \frac{1}{\lambda^4} \left| \frac{(\epsilon - 1)(\epsilon + 2)}{3} \right|^2 \frac{\Delta V^2}{V}, \quad (\text{Gaussian units})$$

(Einstein, 1910; see Jackson, *Classical Electrodynamics*, Sec. 10.2 D). Here  $\epsilon$  is the dielectric constant of the medium,  $\lambda$  is the wavelength of the light, and  $\Delta V^2$  is the mean-square fluctuation  $\Delta V^2 = \langle V^2 \rangle - \langle V \rangle^2$ .

Obtain a general expression for the fluctuation  $\Delta V^2$  in terms of derivatives of  $\mathcal{Y}$ . Show that  $\Delta V^2 = kTV\kappa_T$ , where  $\kappa_T$  is the isothermal compressibility of the medium,  $\kappa_T = -\frac{1}{V} \left( \frac{\partial V}{\partial P} \right)_T$ . (This is another example of the relation between *fluctuations* and the *linear response* of a system.)

- (c) Obtain an expression for  $\kappa_T$  in terms of the critical pressure and volume  $P_c, V_c$ , and the scaled variable  $(V - V_c)/V_c$ , for a van der Waals gas near the critical point on its critical isotherm  $T = T_c$  (see Landau and Lifshitz, Secs. 76, 84). [Hint: expand the expression for  $\partial P/\partial V$  which follows from Eq. (84.5) in a Taylor series in powers of  $(V - V_c)$ , and keep only the first nonzero term. How many derivatives should vanish at the critical point on the  $P$ - $V$  diagram?] Estimate the absorption length  $\ell = \alpha^{-1}$  for blue light ( $\lambda = 4.5 \times 10^{-5}$  cm) in  $\text{CO}_2$  ( $T_c = 304$  K,  $P_c = 72.9$  atm,  $n_c = 8P_c/3kT_c$ ) for  $(V - V_c)/V_c = 10^{-2}$ .  $\epsilon = 1 + 4\pi n\gamma$ , where  $n = N/V$  and  $\gamma$  is the polarizability of the molecule.  $\epsilon = 1.000985$  for  $\text{CO}_2$  at 273 K and 1 atm (but not at  $T_c, P_c$ !). Gaseous and liquid  $\text{CO}_2$  are normally transparent. Comment—relative to your result—on the transparency near the critical point.

### LD 16: Effect of binary collisions on the entropy and energy of a gas

- (a) Derive expressions for the chemical potential, entropy, and total energy of a non-relativistic monatomic gas that take the effects of binary interactions through a potential  $V(r)$  into account to first order in the cluster integral  $b_2$ . Express the results in terms of  $N, T$ , and the number density  $n = N/V$ , and show that the changes from the results for the ideal gas are given by

$$\begin{aligned} \Delta E &= NkT \cdot n\lambda^3 \left( T \frac{db_2}{dT} - \frac{3}{2}b_2 \right), \\ \Delta S &= Nk \cdot n\lambda^3 \left( T \frac{db_2}{dT} - \frac{1}{2}b_2 \right). \end{aligned}$$

[Hints: Start with the cluster expansion for  $\Omega = -kT \ln \mathcal{Z}$ . Solve the equation for  $\mu$  or  $z$  by iteration correct to first order in  $b_2$ . Recall that  $E = -(\partial \ln \mathcal{Z} / \partial \beta)_{V,z}$ . Eliminate  $z$  and introduce  $N$  immediately in the expressions for  $E$  and  $S$ .]

- (b) Write the results for  $\Delta S$  and  $\Delta E$  in terms of integrals involving  $V(r)$ . Show that  $\Delta S \leq 0$  whether  $V(r)$  is attractive or repulsive, *i.e.*, that binary collisions reduce the entropy. Why is this expected? Give a simple interpretation of the result for  $\Delta E$ . [Hint: The inequality  $(1+x)e^{-x} \leq 1$  for all  $x$ ,  $-\infty < x < \infty$  will be useful.]

**LD 17: Interatomic potential for  $^4\text{He}$  from the second virial coefficient.**

The following are measured values (in molar units) of the second virial coefficient  $B(T)$  for helium gas:

$$PV = mN_A kT \left( 1 + \frac{m}{V} B(T) + \dots \right), \quad m = \# \text{ of moles}$$

$T(K)$      $B(T)$  ( $\text{cm}^3/\text{mole}$ )

15	-8.7
20	-2.2
30	3.8
40	6.6
50	8.2
100	11.4
200	12.3
273	12.0
373	11.3

Assume that the interaction of the helium atoms can be described by a Lennard-Jones “6–12 potential”

$$V(r) = -V_0 \left[ 2 \left( \frac{a}{r} \right)^6 - \left( \frac{a}{r} \right)^{12} \right].$$

Approximate the short-distance contributions to  $B(T)$  using a hard core interaction, and calculate the long-distance contribution through order  $(V_0/kT)^2$ . Use the measurements of  $B(T)$  at 30 K and 50 K to determine the potential depth  $V_0$  in temperature units (K) and the location of the minimum ( $r = a$ ). Calculate  $B(T)$  at the temperatures above, and compare your results with the data (give a graph). Why are the deviations at large  $T$  expected? [Caution: Use sufficient accuracy in your numerical calculations.]